

Lecture Wednesday 5 September 2007

Topics Helium part 3: Ground State; Independent electrons; (Variational Method next)

Bound states - energies - and Helium; Classically unstable

In classical physics bound "states" are never periodic & never stable!

Poincaré:

Wavefunction for 2 independent electrons
Ground state of hydrogenlike each one

$$\psi_{1s}^{(z)}(\vec{r}_1) \quad \psi_{1s}^{(z)}(\vec{r}_2) \quad 1s \quad n=1 \quad l=0$$

$$H = T_1 + V_1 + T_2 + V_2 + \frac{e^2}{|\vec{r}_1 - \vec{r}_2|}$$

SPIN MUST BE SINGLET

$$H \left(\psi_{1s}(\vec{r}_1) \phi \psi_{1s}(\vec{r}_2) \right) \rightarrow ?$$

Independent particles - energies

$$(T_1 + V_1 + T_2 + V_2) \psi_a(1) \psi_b(2)$$

hydrogen-like

$$\psi_b(2) (T_1 + V_1) \psi_a(1) + \psi_a(1) (T_2 + V_2) \psi_b(2)$$

$\frac{d^2}{dx^2}$ type

Thus independent electron assumption leads to 2-hydrogenlike Schröd. eq.

$\downarrow E_a \psi_a(1)$ $\downarrow E_b \psi_b(2)$

$$H_0 \psi_a(1) \psi_b(2) = (E_a + E_b) \psi_a(1) \psi_b(2)$$

is the H with neglected $\frac{e^2}{|\vec{r}_1 - \vec{r}_2|}$

The energies are then $\left(\frac{Z^2}{n_a^2} + \frac{Z^2}{n_b^2} \right) E_{1s}(\text{hydrogen})$

Experiment etc

Two-electron systems in units of eV

	H ⁺	He	Li ⁺	Be ⁺⁺	B ⁽³⁺⁾	C ⁽⁴⁺⁾
Z	1	2	3	4	5	6
Ion.pot.[au]	0,757	24,60	75,62	153,95	259,49	392,22
2.Ion.pot	13,600	54,40	122,40	217,60	340,00	489,60
EXP.BindEner	-14,357	-79,00	-198,02	-371,55	-599,49	-881,82
2 E _{1s}	-27,200	-108,80	-244,80	-435,20	-680,00	-979,20
2 E _{1s} + 5/8 Z	-10,200	-74,80	-193,80	-367,20	-595,00	-877,20
Variational	-12,856	-77,46	-196,46	-369,86	-597,66	-879,86
EXP.BindEner	-14,357	-79,00	-198,02	-371,55	-599,49	-881,82

Ion.pot. Ionization potential: The energy to remove the first electron

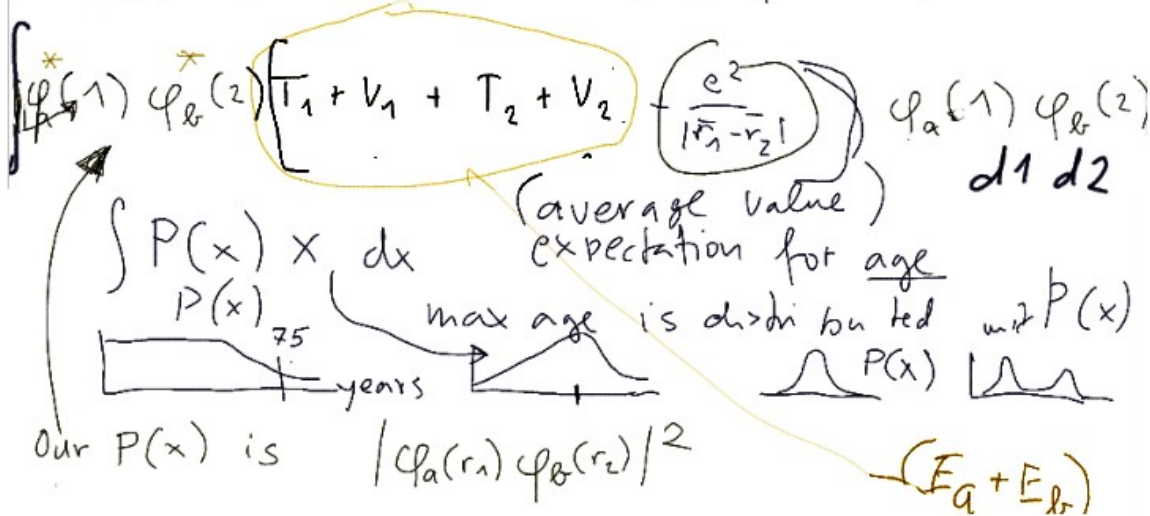
2.Ion.pot Second Ionization potential: The energy to remove the second electron

Evaluate expectation value - What is expectation value ?

Example - age distribution - average value of age

$$\left(T_1 + V_1 + T_2 + V_2 + \frac{e^2}{|\vec{r}_1 - \vec{r}_2|} \right) \psi_a(1) \psi_b(2)$$

Expectation value:
Wavefunction independent;
But CORRECT FULL ENERGY



Generally for any two states a,b - but we specify a→1s b→1s; Multipole expansion

$$\int d^3r_1 d^3r_2 |\varphi_a(r_1)|^2 |\varphi_b(r_2)|^2 \frac{e^2}{|\vec{r}_1 - \vec{r}_2|} \left. \vphantom{\int} \right\} \text{all mixed}$$

$$\varphi_{nlm}(\vec{r}) = R_{nl}(r) Y_{lm}(\hat{r}) \quad \begin{array}{l} \text{angular} \\ \text{variables} \end{array}$$

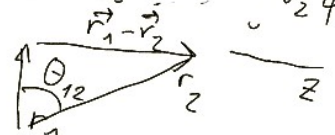
$$d^3r_1 = \underbrace{r_1^2 dr_1}_{\text{only } r} \cdot \underbrace{\sin\theta_1 d\theta_1 d\varphi_1}_{\text{only angles}}$$

$$\frac{1}{|\vec{r}_1 - \vec{r}_2|} = \frac{1}{\sqrt{r_1^2 + r_2^2 - 2r_1 r_2 \cos\theta_{12}}} \quad \text{express } \theta_{12} \text{ in } \begin{array}{l} \theta_1, \varphi_1 \\ \theta_2, \varphi_2 \end{array}$$

General method

$$\frac{1}{|\vec{r}_1 - \vec{r}_2|} = \sum_{L=0}^{\infty} \sum_{M=-L}^L \frac{4\pi}{2L+1} \left(\frac{r_2^L}{r_1^{L+1}} \right) Y_{LM}(\theta_1, \varphi_1) Y_{LM}(\theta_2, \varphi_2)$$

(multipole expansion of Coulomb term)



How to perform the 6-dim integration: two angular (4 dimensions) using Y - spherical harmonics

$$\int d^3r_1 d^3r_2 |\varphi_a(r_1)|^2 |\varphi_b(r_2)|^2 \frac{e^2}{|\vec{r}_1 - \vec{r}_2|} \left. \vphantom{\int} \right\} \text{all mixed}$$

$$\varphi_{nlm}(\vec{r}) = R_{nl}(r) Y_{lm}(\hat{r}) \quad \begin{array}{l} \text{angular} \\ \text{variables} \end{array}$$

$$d^3r_1 = \underbrace{r_1^2 dr_1}_{\text{only } r} \cdot \underbrace{\sin\theta_1 d\theta_1 d\varphi_1}_{\text{only angles}}$$

$$\varphi_{1s} \rightarrow C e^{-\frac{Zr}{a_0}} \cdot Y_{00}(\theta, \varphi) \quad Y_{00} = \frac{1}{\sqrt{4\pi}} \quad \text{constant}$$

$$\int Y_{lm}^*(\theta, \varphi) Y_{l'm'}(\theta, \varphi) d\Omega = \delta_{l,l'} \delta_{m,m'}$$

$$\int Y_{lm} d\Omega = 0 \quad \text{for } l \neq 0 \text{ because } \int Y_{lm} Y_{00}^* \rightarrow 0$$

$$\delta_{ab} = \begin{cases} 0 & a \neq b \\ 1 & \text{otherwise} \end{cases}$$

Orthogonality of spherical harmonics

Gives that the integral of one Y is zero (unless L=0....)

The evaluation reduces to a two-dimensional integral over r_1 and r_2 (with $r_>$)

$$\iiint R_{00}^* Y_{00}^* \left(\sum_{\substack{L=0 \\ M=0}}^{\infty} \frac{(4\pi)}{2L+1} \left(\frac{r_{<}}{r_{>}} \right)^L Y_{LM}^*(\vartheta_1, \varphi_1) Y_{LM}(\vartheta_2, \varphi_2) \right) R_{00} Y_{00}$$

for general LM the angular integrals

GIVE ZERO! $Y_{00} = \frac{1}{\sqrt{4\pi}}$
only $L=0 M=0$ can contribute

all angular integrals combined with give \uparrow

$$\int r_1^2 dr_1 \int r_2^2 dr_2 \frac{1}{r_{>}} r_{>}^2$$

$$\frac{1}{|\vec{r}_1 - \vec{r}_2|} = \sum_{L=0}^{\infty} \sum_{M=-L}^L \frac{(4\pi)}{2L+1} \left(\frac{r_{<}}{r_{>}} \right)^L Y_{LM}^*(\vartheta_1, \varphi_1) Y_{LM}(\vartheta_2, \varphi_2)$$

We must evaluate the integrals

$$R_{1,0}(r) = 2 \cdot \left(\frac{Z}{a_0} \right)^{\frac{3}{2}} \cdot e^{-\frac{Zr}{a_0}} = R_{1,0}^*(r)$$

$$\begin{aligned} & \int_0^\infty \int_0^\infty r_1^2 \cdot r_2^2 \cdot R_{1,0}(r_1)^2 \cdot R_{1,0}(r_2)^2 \frac{e^2}{r_{>}} dr_1 dr_2 \\ &= \int_0^\infty \int_0^\infty 2^4 \left(\frac{Z}{a_0} \right)^6 e^{-\frac{2Z}{a_0}(r_1+r_2)} r_1^2 \cdot r_2^2 \frac{e^2}{r_{>}} dr_1 dr_2 \\ &= 2^4 \left(\frac{Z}{a_0} \right)^6 \cdot e^2 \int_0^\infty \int_0^\infty e^{-\frac{2Z}{a_0}(r_1+r_2)} r_1^2 \cdot r_2^2 \frac{1}{r_{>}} dr_1 dr_2 \end{aligned}$$

The $r_{>}$ must be replaced which leads to two integrations

$$\begin{aligned} & \int_0^\infty \left(\int_0^{r_1} e^{-r_1-r_2} r_1 r_2^2 dr_2 \right) dr_1 + \int_0^\infty \left(\int_{r_1}^\infty e^{-r_1-r_2} r_1^2 r_2 dr_2 \right) dr_1 \\ &= \int_0^\infty r_1 e^{-r_1} \underbrace{\int_0^{r_1} r_2^2 e^{-r_2} dr_2}_{intB} dr_1 + \int_0^\infty r_1^2 e^{-r_1} \underbrace{\int_{r_1}^\infty e^{-r_2} r_2 dr_2}_{intC} dr_1 \end{aligned}$$

And the result is

$$\langle |V_{12}| \rangle \rightarrow \int_0^\infty \int_0^\infty \dots = \frac{5}{8} \frac{Ze^2}{a_0}$$

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$$\langle |V_{12}| \rangle \rightarrow \int_0^\infty \int_0^\infty \dots = \frac{5}{8} \frac{Ze^2}{a_0}$$

The multipole expansion we used:

$$\frac{1}{|\vec{r}_1 - \vec{r}_2|} = \sum_{LM} \frac{4\pi}{2L+1} \frac{r_{<}^L}{r_{>}^{L+1}} Y_{LM}^*(\hat{r}_1) Y_{LM}(\hat{r}_2)$$

$$r_{<} = r_1, \quad r_{>} = r_2 \quad \text{for} \quad |\vec{r}_1| < |\vec{r}_2|$$

$$r_{<} = r_2, \quad r_{>} = r_1 \quad \text{for} \quad |\vec{r}_1| > |\vec{r}_2|$$