PHYS261 Atomic Physics and Physical Optics

Lecture Tuesday 2. September 2008

Topics: 1. Selection Rules, Parity - hydrogen-like levels2. Helium Atom - starting the topic

Comment: Selection rules, transition probability The theory of eletromagnetic transitions will be treated in detail later - in "Atoms and light" section





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"Simplified picture of a 'transition'

the transition probability – transition amplitudes – wavefunction integrals – are (non) zero

Amplitude:
$$\langle \gamma_{f}| H_{I} | \gamma_{i} \rangle$$

Probabilities of transition $|amplitude|^{2}$
 $\langle \gamma_{f}| H_{I}| \gamma_{i} \rangle = \int d^{3}\vec{r} \gamma_{f}^{*}(\vec{r}) H_{i}(\vec{r}) \gamma_{i}(\vec{r})$
These must be different from O
 $H_{i} \propto \vec{r}$ (dipole)

Simplified picture of a 'transition'

the transition probability – transition amplitudes – wavefunction integrals – are (non) zero

The atom works as ' antenna ' - the dipole antenna is the simplest most efficient

 $m_l = 0 \longrightarrow$ n, l=0=0 const in angle dep on absolute distance i'n angle J Ymjo F Vhz, o d'F \sim 0 $\int Y_{n_1\ell=1} \stackrel{*}{r} \frac{1}{\gamma_{n_2\ell}} \frac{1}{\gamma_{n_2\ell}$ *‡* 0 ~ S(r' · r') Yhod I dr /to each minus contribution /2 outy >0 K=a+ib ~= a-ib △L = ± 1 gives houzero $\Delta L = l_f - l'_i$



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Homework on parity:

A mixed parity function is just any function | f(x) | not equal to | f(-x) |

Find the decomposition of such general function f(x)into a positive parity p(x) and negative parity n(x), i.e. express p(x) and n(x) in terms of f(x)

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Helium
2 electrons
                    2 charges -1
nucleus: 2 protons + 2 neutrons
                 total charge +2
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$$H(\vec{r_{1}},\vec{r_{2}}) \quad \Psi(\vec{r_{1}},\vec{r_{2}}) = E \quad \Psi(\vec{r_{1}},\vec{r_{2}})$$
Start with independent electrons
(forget about $\vec{r_{2}}$)
Probability density
$$P(\vec{r_{1}},\vec{r_{2}})|^{2} \qquad (forget about $\vec{r_{2}}$)
$$P(\vec{r_{1}},\vec{r_{2}})|^{2} \qquad e_{q} \text{ at } r_{q} \quad \text{"at the } \cdot \cdot \cdot \cdot e_{q} \text{ at } r_{2} \quad \text{same hue"}$$

$$(correlation) \quad e_{2} \text{ at } r_{2} \quad \text{same hue"}$$

$$two independent events \qquad product of probabilities$$

$$\Psi(\vec{r_{1}},\vec{r_{2}}) = \Psi_{1}(\vec{r_{1}}) \cdot \Psi_{2}(\vec{r_{2}})$$

$$P(\vec{r_{1}},\vec{r_{2}}) = \Psi_{1}(\vec{r_{1}}) \cdot \Phi_{2}(\vec{r_{2}})$$

$$Independent probability densities$$

$$\Psi(\vec{r_{1}},\vec{r_{2}}) = \Psi_{1}(\vec{r_{1}}) \cdot \Phi_{2}(\vec{r_{2}})$$

$$P(\vec{r_{1}},\vec{r_{2}}) = \Psi_{1}(\vec{r_{1}}) \cdot \Phi_{2}(\vec{r_{2}})$$$$

$$\begin{bmatrix} T_1 + V_1 + T_2 + V_2 \end{bmatrix} \Psi_4(r_1) \Psi_2(r_2) = E \Lambda \Psi_4(r_1) \Psi_2(r_2)$$

$$\begin{bmatrix} T_1 + V_1 \end{pmatrix} \Psi_4(r_1) \end{bmatrix} \Lambda \Psi_2(r_2) + \begin{bmatrix} T_2 + V_2 \end{bmatrix} \Psi_4(r_1) \end{bmatrix} \Psi_4(r_1)$$

$$= E \Lambda \Psi_1(r_1) \Psi_1(r_1)$$

$$= E \Lambda \Psi_1(r_1) \Psi_2(r_2)$$

$$= E \Lambda \Psi_1(r_1) \Psi_2(r_2)$$
Separation of variables $\Psi_2(r_2)$

$$= E = E_1 + E_2$$

$$\begin{bmatrix} T_1 + V_1 \end{pmatrix} \Psi_4(r_1) = E_1 \Psi_1(r_1) \qquad (and He) \\ Same eq. \\ for 2 \end{bmatrix}$$

Product wavefunction -, still not the whole story Because of the <u>spin</u> (Pauli) So called Pauli Principle Aufban Prinzip (Build up) In each state only 1 electron Spin came from another "story"

The story of 1924-1930 development of Quantum Mechanics Atomic Theory

and Quantum Chemistry is fascinating!!

There were both Bohr and Heisenberg, but also Pauli, Heitler and London, and also Pauling