

This note was written in "realtime" by Ingjald Pilskog. Edited by L. K.

## Third lecture 29.08.06

*Atomic world construction*

$a_0$  - radius, length

$$\Delta x \approx a_0 \quad \Delta k \approx 1/a_0$$

$(\hbar \Delta k = \Delta p)$  wavenumber

$$\Delta E \approx \frac{(\Delta p)^2}{2m} = \frac{\hbar^2}{2ma_0^2} \approx \frac{\hbar^2}{ma_0^2} \quad (\text{kinetic energy } T_0. \quad \text{Potential energy } V_0 \mid - \frac{e^2}{a_0})$$

**Atomic unit of length**

$$a_0 = 0,529 \text{ \AA} = \frac{\hbar^2}{e^2 m} \text{ Length unit}$$

*Atomic unit of energy*

$$\text{Energy unit } \frac{e^2}{a_0} = \frac{\hbar}{ma_0^2} = 27,2 \text{ eV}$$

$\langle T_0 \rangle = \langle V_0 \rangle$  should be of the same order.

*Atomic unit of velocity*

$$p_0 = \hbar k_0 = \frac{\hbar}{a_0} \quad v_0 = \frac{p_0}{m} = \frac{\hbar}{m a_0} = \frac{me^2 \hbar}{\hbar^2 m} = \frac{e^2}{\hbar} \quad \frac{v_0}{c} = \frac{e^2}{\hbar c} = \alpha = \frac{1}{137}$$

*Atomic unit of time*

$$t_0 = \frac{a_0}{v_0} = \left(\frac{a_0}{e^2}\right) \hbar = \frac{\hbar}{E_0}$$

Alternative postulate  $t_0 = \frac{\hbar}{E_0}$

(Statement a.u.  $\longleftrightarrow e = m\hbar = 1$  is useless)

$$\hbar = 0,66 \cdot 10^{-15} \text{ eVs} \quad t_0 = \frac{0,66 \cdot 10^{-15}}{27,2} \text{ s} = 0,24 \cdot 10^{-15}$$

( $2\pi$  for angular freq.)

$\nu$  frequency  $\omega$  - angular frequency

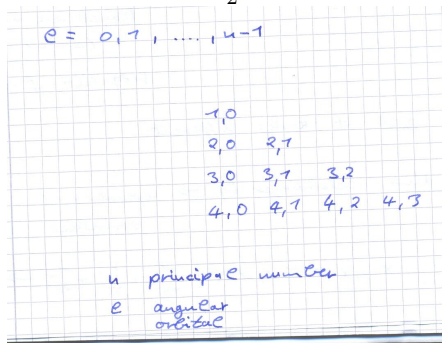
$k_0 T$  is the "physical temperature".

Room temperature is thus  $\frac{1}{40}$  eV or 25 meV

Atomic unit of energy  $\rightarrow$  VERY HOT

### More about bound states in H

Ground state  $-\frac{1}{2}$  a.u.



States characterized by  $n, l$  ( $m$  - magnetic)

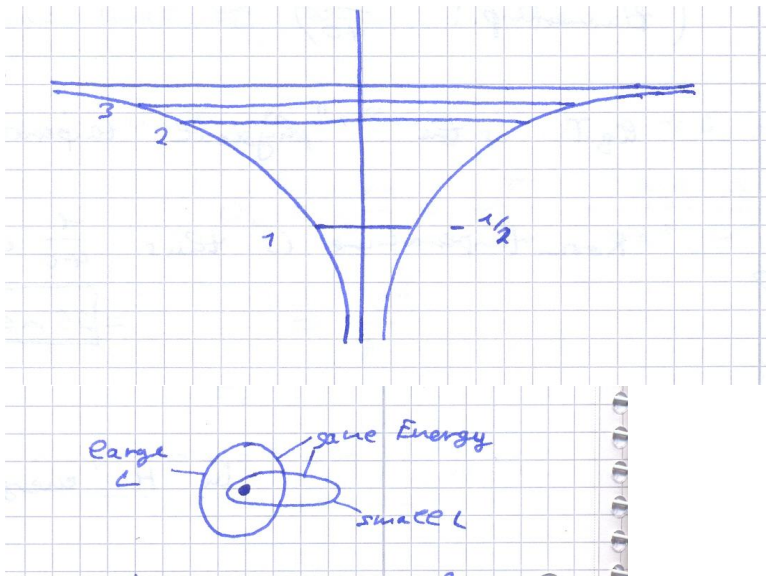
In H energies given by  $\frac{1}{n^2}(-\frac{1}{2})$  a.u.

"Ladi says its nonsens to talk about  $m$  as a quantum number", wrote Ingjald. He did not write: "unless we have magnetic field on".

Angular momentum

$$L = \omega \mathcal{I} \quad T_{rot} = \frac{1}{2} \frac{L^2}{\mathcal{I}} \quad T_{rot} = \frac{1}{2} \mathcal{I} \omega^2 \quad \mathcal{I} \text{ is the moment of inertia}$$

$E = T + V$  is negativ ....  $T$  is kinetic energy



$I_n$  QM it looks different

3-dim Schr.Eq. → Separation of variables

$x, y, z, \rightarrow r, \nu, \varphi$

$$\frac{\delta^2}{\delta x^2} + \frac{\delta^2}{\delta y^2} + \frac{\delta^2}{\delta z^2} \rightarrow T_r + \frac{L^2(\theta, \varphi)}{r^2}$$

$$T_r \rightarrow \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r}$$

$L^2$  is ugly (but can be made very elegant)

This is generally used in many fields.

$$\frac{L^2(\theta, \varphi)}{r^2} \rightarrow \frac{1}{r^2} \frac{1}{\sin\theta} \frac{\partial}{\partial \theta} \left( \sin\theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2 \psi}{\partial \phi^2}$$

Solutions of the above are  $Y_l^m(\theta, \phi)$ , known as the spherical harmonics. (The constants were not taken care of)

From the web ("stolen latex code")

<http://vergil.chemistry.gatech.edu/notes/quantrev/node25.html>

$$-\frac{\hbar^2}{2\mu} \left[ \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial \theta} \left( \sin\theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2 \psi}{\partial \phi^2} \right] - \frac{e^2}{4\pi\epsilon_0 r} \psi(r, \theta, \phi) = E\psi(r, \theta, \phi) \quad (1)$$

Exercise: Look on the separation of variables and how it's done.

( $l=0$ ) s-states,

( $l=1$ ) p-states,

( $l=2$ ) d-states,

( $l=3$ ) f-states, ... there are more, but it is not relevant in typical atoms